

Problem 1.1

A particle of mass  $m$  moves on a circular orbit of radius  $r$  about a central point to which it is attracted by force  $F = \alpha r$  ( $\alpha = \text{const}$ ). Use classical mechanics augmented with the second Bohr's postulate ( $pr = n\hbar$ ) to find the allowed values of:

- (i) orbit radius,
- (ii) particle velocity,
- (iii) kinetic energy, and
- (iv) potential energy.

Compare the last two results with each other and with the expression  $\hbar\omega$ , where  $\omega$  is the frequency of 1D oscillations of the same system. Give a (very brief) discussion of your findings.

*Solution:*

Combining the classical equation of motion  $mv^2/r = \alpha r$  with the Bohr's postulate, we get:

- (i)  $r_n = n^{1/2}x_0$ , where  $x_0 \equiv (\hbar/m\omega)^{1/2}$ , and  $\omega = (\alpha/m)^{1/2}$ .
- (ii)  $v_n = \omega r_n = n^{1/2}\omega x_0$ ,
- (iii)  $T_n = mv_n^2/2 = n\hbar\omega/2$ ,
- (iv)  $U_n = \alpha r_n^2/2 = n\hbar\omega/2$ .

The total energy of the system,  $E_n = T_n + U_n = n\hbar\omega$ . This means that the frequency of radiation at quantum transitions between two adjacent energy levels ( $\Delta E = E_{n+1} - E_n$ ) is exactly equal to the classical frequency of this system (a "harmonic oscillator").

Problem 1.2

In class, we have calculated the Fourier amplitude distribution  $a(k) = C \times \exp(-\frac{k^2}{2\sigma_k^2})$

corresponding to the Gaussian wave packet  $\psi(x) = \exp(-\frac{x^2}{2\sigma_x^2})$ , with  $\sigma_x \sigma_k = 1$ .

- (i) Calculate the constant factor  $C$  in the expression for  $a(k)$ .
- (ii) Find  $a(k)$  for the wave packet  $\psi(x,0) = \exp(ik_0x) \exp(-\frac{x^2}{2\sigma_x^2})$ , where  $k_0 = \text{const}$ .

*Solution:*

- (i) Direct integration, using the method discussed in class and the table integral

$$\int_{-\infty}^{+\infty} \exp(-\xi^2) d\xi = \sqrt{\pi}, \text{ gives } C = \frac{\sigma_x}{\sqrt{2\pi}} = \frac{1}{\sqrt{2\pi}\sigma_k}.$$

- (ii) Since  $a(k) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \psi(x,0) \exp(-kx) dx$ , we see that the multiplication of  $\psi(x,0)$  by

$$\exp(ik_0x) \text{ just shifts the argument } k \text{ by } k_0, \text{ i.e. } a(k) = C \times \exp[-\frac{(k - k_0)^2}{2\sigma_k^2}].$$

### Problem 1.3

Calculate the wave packet  $\psi(x,0)$  corresponding to the distribution  $a(k) = \frac{1}{(k - k_0)^2 + \kappa^2}$ , where  $k_0$  and  $\kappa$  are constants.

*Solution:*

Using the table integral  $\int_{-\infty}^{+\infty} \frac{\cos ax}{a^2 + 1} dx = \pi \exp(-|a|)$ , we get  $\psi(x,0) = \frac{\pi}{|\kappa|} \exp(ik_0x - |\kappa x|)$ .

### Problem 1.4

Find the phase and group velocities for waves with the following dispersion laws:

(i)  $\lambda = \frac{c}{\sqrt{\nu^2 - \nu_0^2}}$  (e.g., EM waves in plasma or inside a hollow waveguide),

(ii)  $\nu = \sqrt{\frac{g}{2\pi\lambda}}$  (gravity waves on deep water).

$c$ ,  $\nu_0$  and  $g$  (and of course  $\pi \cong 3.14$ ) are constants.

*Solution:*

In both cases, we should first rewrite the dispersion relation in the form  $\omega = \omega(k)$  (using the relations  $\nu = \omega/2\pi$ ,  $\lambda = 2\pi/k$ ), and then calculate  $v_p = \omega/k$  and  $v_g = \partial\omega/\partial k$ . This procedure yields:

(i)  $\omega = \sqrt{\omega_0^2 + k^2 c^2}$ ,  $v_p = \frac{\sqrt{\omega_0^2 + k^2 c^2}}{k}$ ,  $v_g = \frac{kc^2}{\sqrt{\omega_0^2 + k^2 c^2}}$ , so that  $v_p v_g = c^2$ ;

(ii)  $\omega = \sqrt{gk}$ ,  $v_p = \sqrt{\frac{g}{k}}$ ,  $v_g = \frac{1}{2} \sqrt{\frac{g}{k}}$ , so that  $v_g = v_p/2$ .